



Basics

- Formulation
- Basis functions
- Stiffness matrix

Poisson's equation

- Regular grid
- Boundary conditions
- Irregular grid

Numerical Examples





Let us start with a simple linear system of equations



and observe that we can generally multiply both sides of this equation with y without changing its solution. Note that x, y and b are vectors and A is a matrix.

$$\rightarrow$$
 yAx= yb $y \in \Re^n$

We first look at Poisson's equation

$$-\Delta u(x) = f(x)$$

where u is a scalar field, f is a source term and in 1-D

$$\Delta = \nabla^2 = \frac{\partial^2}{\partial x^2}$$



We now multiply this equation with an arbitrary function v(x), (dropping the explicit space dependence)

$$-\Delta uv = fv$$

... and integrate this equation over the whole domain. For reasons of simplicity we define our physical domain D in the interval [0, 1].

$$-\int_{D} \Delta uv = \int_{D} fv$$
$$-\int_{0}^{1} \Delta uv dx = \int_{0}^{1} fv dx$$

Das Reh springt hoch, das Reh springt weit, warum auch nicht, es hat ja Zeit.

... why are we doing this? ... be patient ...



As we are aiming to find a numerical solution to our problem it is clear we have to discretize the problem somehow. In FE problems - similar to FD - the functional values are known at a discrete set of points.

Domain D

The key idea in FE analysis is to approximate all functions in terms of basis functions $\phi,$ so that

$$u \approx \widetilde{u} = \sum_{i=1}^{N} c_i \varphi_i$$





$$u \approx \widetilde{u} = \sum_{i=1}^{N} c_i \varphi_i$$

where N is the number nodes in our physical domain and c_i are real constants.

With an appropriate choice of basis functions φ_i , the coefficients c_i are equivalent to the actual function values at node point i. This - of course - means, that φ_i =1 at node i and 0 at all other nodes ...

Doesn't that ring a bell?

Before we look at the basis functions, let us ...





... partially integrate the left-hand-side of our equation ...

$$-\int_{0}^{1} \Delta uv dx = \int_{0}^{1} fv dx$$
$$-\int_{0}^{1} (\nabla \bullet \nabla u) v dx = \boxed{[\nabla uv]_{0}^{1}} + \int_{0}^{1} \nabla v \nabla u dx$$

we assume for now that the derivatives of u at the boundaries vanish so that for our particular problem

$$-\int_{0}^{1} (\nabla \bullet \nabla u) v dx = \int_{0}^{1} \nabla v \nabla u dx$$





... so that we arrive at ...

$$\int_{0}^{1} \nabla u \nabla v dx = \int_{0}^{1} f v dx$$

... with u being the unknown. This is also true for our approximate numerical system

$$\int_{0}^{1} \nabla \widetilde{u} \nabla v dx = \int_{0}^{1} f v dx$$

... where ...

$$\widetilde{u} = \sum_{i=1}^{N} c_i \varphi_i$$

was our choice of approximating u using basis functions.



$$\int_{0}^{1} \nabla \widetilde{u} \nabla v dx = \int_{0}^{1} f v dx$$

... remember that v was an arbitrary real function ... if this is true for an arbitrary function it is also true if

$$v = \varphi_j$$

... so any of the basis functions previously defined ...

... now let's put everything together ...



... leading to ...





... the coefficients c_k are constants so that for one particular function ϕ_k this system looks like ...





The solution

... with the even less surprising solution

$$b_i = \left(A_{ik}^T\right)^{-1} g_k$$

remember that while the b_i's are really the coefficients of the basis functions these are the actual function values at node points i as well because of our particular choice of basis functions.

This become clear further on ...



The basis functions

9

8

7

6

5

4

3

2

1



we are looking for functions ϕ_i with the following property

... otherwise we are free to choose any function ...

The simplest choice are of course linear functions:

+ grid nodes

blue lines - basis functions ϕ_i







To assemble the stiffness matrix we need the gradient (red) of the basis functions (blue)







Knowing the particular form of the basis functions we can now calculate the elements of matrix A_{ij} and vector g_i

$$\sum_{i=1}^{n} c_{i} \int_{0}^{1} \nabla \varphi_{i} \nabla \varphi_{k} dx = \int_{0}^{1} f \varphi_{k} dx \qquad \Rightarrow \qquad b_{i} A_{ik} = g_{k}$$
$$A_{ik} = \int_{0}^{1} \nabla \varphi_{i} \nabla \varphi_{k} dx \qquad \qquad g_{k} = \int_{0}^{1} f \varphi_{k} dx$$

Note that φ_i are continuous functions defined in the interval [0,1], e.g.

$$\varphi_{i}(x) = \begin{cases} \frac{x - x_{i-1}}{x_{i} - x_{i-1}} & \text{for } x_{i-1} < x \le x_{i} \\ \frac{x_{i+1} - x}{x_{i+1} - x_{i}} & \text{for } x_{i} < x < x_{i+1} \\ 0 & \text{elsewhere} \end{cases}$$
Let us - for now - assume regular grid ... then

now - assume a





$$\varphi_{i}(x) = \begin{cases} \frac{x - x_{i-1}}{x_{i} - x_{i-1}} & \text{for } x_{i-1} < x \le x_{i} \\ \frac{x_{i+1} - x}{x_{i+1} - x_{i}} & \text{for } x_{i} < x < x_{i+1} \\ 0 & \text{elsewhere} \end{cases} \Rightarrow \quad \varphi_{i}(\tilde{x}) = \begin{cases} \frac{\tilde{x}}{dx} + 1 & \text{for } -dx < \tilde{x} \le 0 \\ 1 - \frac{\tilde{x}}{dx} & \text{for } 0 < \tilde{x} < dx \\ 0 & \text{elsewhere} \end{cases}$$

... where we have used ...



$$\widetilde{x} = x - x_i$$
$$dx = x_i - x_{i-1}$$



$$\nabla \varphi_i(\widetilde{x}) = \begin{cases} 1/dx & \text{for } -dx < \widetilde{x} \le 0\\ -1/dx & \text{for } 0 < \widetilde{x} < dx\\ 0 & \text{elsewhere} \end{cases}$$

$$\begin{aligned} \widetilde{x} &= x - x_i \\ dx &= x_i - x_{i-1} \end{aligned}$$







$$A_{ik} = \int_{0}^{1} \nabla \varphi_{i} \nabla \varphi_{k} dx$$

... we have to distinguish various cases ... e.g. ...

$$A_{11} = \int_{0}^{1} \nabla \varphi_{1} \nabla \varphi_{1} dx = \int_{x_{1}}^{x_{1}+dx} \nabla \varphi_{1} \nabla \varphi_{1} dx = \int_{x_{1}}^{x_{1}+dx} \frac{-1}{dx} \frac{-1}{dx} \frac{-1}{dx} \frac{-1}{dx} dx = \frac{1}{dx^{2}} \int_{0}^{dx} dx = \frac{1}{dx}$$
$$A_{22} = \int_{0}^{1} \nabla \varphi_{2} \nabla \varphi_{2} dx = \int_{x_{2}-dx}^{x_{2}} \nabla \varphi_{2} \nabla \varphi_{2} dx + \int_{x_{2}}^{x_{2}+dx} \nabla \varphi_{2} \nabla \varphi_{2} dx$$
$$= \frac{1}{dx^{2}} \int_{-dx}^{0} dx + \frac{1}{dx^{2}} \int_{0}^{dx} dx = \frac{2}{dx}$$



$$A_{ik} = \int_{0}^{1} \nabla \varphi_i \nabla \varphi_k dx$$



... and ...

$$A_{12} = \int_{0}^{1} \nabla \varphi_{1} \nabla \varphi_{2} dx = \int_{x_{1}}^{x_{1}+dx} \nabla \varphi_{1} \nabla \varphi_{2} dx = \int_{x_{1}}^{x_{1}+dx} \frac{-1}{dx} \frac{1}{dx} \frac{1}{dx} dx$$
$$= \frac{-1}{dx^{2}} \int_{0}^{dx} dx = \frac{-1}{dx}$$
$$A_{21} = A_{12}$$

... so that finally the stiffness matrix looks like ...



Stiffness matix - elements

$$A_{ik} = \int_{0}^{1} \nabla \varphi_i \nabla \varphi_k dx$$



$$A_{ij} = \frac{1}{dx} \begin{pmatrix} 1 & -1 & & \\ -1 & 2 & -1 & & \\ & \ddots & & \\ & & -1 & 2 & -1 \\ & & & -1 & 1 \end{pmatrix}$$

... so far we have ignored sources and boundary conditions ...





... let us start restating the problem ...

$$-\Delta u(x) = f(x)$$

... which we turned into the following formulation ...

$$\sum_{i=1}^{n} c_i \int_{0}^{1} \nabla \varphi_i \nabla \varphi_k dx = \int_{0}^{1} f \varphi_k dx$$

... assuming ...

$$\widetilde{u} = \sum_{i=1}^{N} c_i \varphi_i \qquad \text{ with b.c. } \quad \widetilde{u} = \sum_{i=2}^{N-1} c_i \varphi_i + u(0)\varphi_1 + u(1)\varphi_N$$

where u(0) and u(1) are the values at the boundaries of the domain [0,1]. How is this incorporated into the algorithm?



$$\sum_{i=1}^{n} c_i \int_{0}^{1} \nabla \varphi_i \nabla \varphi_k dx = \int_{0}^{1} f \varphi_k dx$$

$$-\Delta u(x) = f(x)$$

... which we turned into the following formulation ...



... the system *feels* the boundary conditions through the (modified) source term



Numerical example - regular grid



$-\Delta u(x) = f(x)$

```
Domain: [0,1]; nx=100;
dx=1/(nx-1); f(x)=\delta(1/2)
Boundary conditions:
u(0)=u(1)=0
```

Matlab FD code

```
f(nx/2)=1/dx;
for it = 1:nit,
uold=u;
du=(csh(u,1)+csh(u,-1));
u=.5*( f*dx^2 + du );
u(1)=0;
u(nx)=0;
end
```

Matlab FEM code

```
% source term
s=(1:nx)*0; s(nx/2)=1.;
% boundary left u_1 int{ nabla phi_1 nabla phij }
111 = 0;
        s(1) = 0;
% boundary right u_nx int{ nabla phi_nx nabla phij }
unx=0; s(nx)=0;
% assemble matrix Aij
A=zeros(nx);
for i=2:nx-1,
   for j=2:nx-1,
      if i==j,
         A(i,j)=2/dx;
      elseif j==i+1
         A(i,j) = -1/dx;
      elseif j==i-1
         A(i,j) = -1/dx;
      else
         A(i, j) = 0;
      end
   end
end
fem(2:nx-1)=inv(A(2:nx-1,2:nx-1))*s(2:nx-1)';
fem(1)=u1;
fem(nx)=unx;
```



$$-\Delta u(x) = f(x)$$





$$-\Delta u(x) = f(x)$$











$$A_{12} = \int_{0}^{1} \nabla \varphi_{1} \nabla \varphi_{2} dx = \int_{x_{1}}^{x_{1}+h_{1}} \nabla \varphi_{1} \nabla \varphi_{2} dx = \int_{x_{1}}^{x_{1}+h_{1}} \frac{-1}{h_{1}} \frac{1}{h_{1}} \frac{1}{h_{1}} dx$$

 A_{ii}





Stiffness matrix A

$$-\Delta u(x) = f(x)$$

Domain: [0,1]; nx=100; dx=1/(nx-1); $f(x)=\delta(1/2)$ Boundary conditions: u(0)=u0; u(1)=u1

i=1	2	3	4	5	6	7
+	+	+	+	+	+	+
	\mathbf{h}_1	h ₂	h ₃	h ₄	h ₅	h ₆

```
for i=2:nx-1,
   for j=2:nx-1,
      if i==j,
         A(i,j)=1/h(i-1)+1/h(i);
      elseif i==j+1
         A(i, i) = -1/h(i-1);
      elseif i+1==j
         A(i,j) = -1/h(i);
      else
         A(i,j)=0;
      end
   end
end
```





FEM on Chebyshev grid





In finite element analysis we approximate a function defined in a Domain D with a set of orthogonal basis functions with coefficients corresponding to the functional values at some node points.

The solution for the values at the nodes for some partial differential equations can be obtained by solving a linear system of equations involving the inversion of (sometimes sparse) matrices.

Boundary conditions are inherently satisfied with this formulation which is one of the advantages compared to finite differences.