



Motivation

Specific methods:

- Finite differences
- High-order FD methods
- Pseudospectral methods
- Finite elements
- Finite volumes

Applications

- Wave propagation
- Rupture problems
- Volcano seismology
- Global wave propagation
- Earthquake scenarios





















Why numerical methods









Spatial Scales and Memory

(back of the envelope)



Highest frequency: Shortest wavelength: Shortest wavelength: Grid points per wavelength: Grid spacing: Grid spacing: 0.1 Hz 20 km (crust) 50 km (mantle) 5 2000 m (crust) 5000 m (mantle)





MAM mary management

Required grid points: O(10⁹) Required memory: O(100 GBytes)







$$\partial_{t}^{2} \mathbf{p} = \mathbf{c}^{2} \Delta \mathbf{p} + \mathbf{s}$$
$$\Delta = (\partial_{x}^{2} + \partial_{y}^{2} + \partial_{z}^{2})$$

Ρ	pressure
С	acoustic wave speed
S	sources

The acoustic wave equation

- seismology
- acoustics
- oceanography
- meteorology

 $\partial_t C = k \Delta C - \mathbf{v} \bullet \nabla C - RC + p$

- C tracer concentration
- k diffusivity
- v flow velocity
- R reactivity
 - sources

D

Diffusion, advection, Reaction

- geodynamics
- oceanography
- meteorology
- geochemistry
- sedimentology
- geophysical fluid dynamics



Numerical methods: properties



Finite differences	 time-dependent PDEs seismic wave propagation geophysical fluid dynamics Maxwell's equations Ground penetrating radar robust, simple concept, easy to parallelize, regular grids, explicit method
Finite elements	 static and time-dependent PDEs seismic wave propagation geophysical fluid dynamics all problems > implicit approach, matrix inversion, well founded, irregular grids, more complex algorithms, engineering problems
Finite volumes	 time-dependent PDEs seismic wave propagation mainly fluid dynamics robust, simple concept, <u>irregular grids</u>, explicit method



Other Numerical methods:









Common definitions of the derivative of f(x):

$$\partial_{x} f = \lim_{dx \to 0} \frac{f(x + dx) - f(x)}{dx}$$

$$\partial_x f = \lim_{dx \to 0} \frac{f(x) - f(x - dx)}{dx}$$

$$\partial_x f = \lim_{dx \to 0} \frac{f(x+dx) - f(x-dx)}{2dx}$$

These are all correct definitions in the limit dx->0.

But we want dx to remain **FINITE**



What is a finite difference?



The equivalent *approximations* of the derivatives are:

$$\partial_x f^+ \approx \frac{f(x+dx) - f(x)}{dx}$$

forward difference

$$\partial_x f^- \approx \frac{f(x) - f(x - dx)}{dx}$$

backward difference

$$\partial_x f \approx \frac{f(x+dx) - f(x-dx)}{2dx}$$

centered difference





Our first FD algorithm (ac1d.m) !



$$\partial_{t}^{2} \mathbf{p} = \mathbf{c}^{2} \Delta \mathbf{p} + \mathbf{s}$$
$$\Delta = (\partial_{x}^{2} + \partial_{y}^{2} + \partial_{z}^{2})$$

Ρ	pressure
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Problem: Solve the 1D acoustic wave equation using the finite Difference method.

Solution:

$$p(t + dt) = \frac{c^2 dt^2}{dx^2} \left[p(x + dx) - 2 p(x) + p(x - dx) \right] + 2 p(t) - p(t - dt) + sdt^2$$





$$p(t + dt) = \frac{c^2 dt^2}{dx^2} \left[p(x + dx) - 2 p(x) + p(x - dx) \right] + 2 p(t) - p(t - dt) + sdt^2$$

Stability: Careful analysis using harmonic functions shows that a stable numerical calculation is subject to special conditions (conditional stability). This holds for many numerical problems.

$$\mathbf{C}\frac{\mathbf{dt}}{\mathbf{dx}} \le \varepsilon \approx 1$$





$$p(t + dt) = \frac{c^2 dt^2}{dx^2} \Big[p(x + dx) - 2 p(x) + p(x - dx) \Big] + 2 p(t) - p(t - dt) + sdt^2$$



Dispersion: The numerical approximation has artificial dispersion, in other words, the wave speed becomes frequency dependent. You have to find a frequency bandwidth where this effect is small. The solution is to use a sufficient number of grid points per wavelength.



Our first FD code!



$$p(t + dt) = \frac{c^2 dt^2}{dx^2} \left[p(x + dx) - 2 p(x) + p(x - dx) \right] + 2 p(t) - p(t - dt) + sdt^2$$

% Time stepping

```
for i=1:nt,
```

```
% FD
```

```
disp(sprintf(' Time step : %i',i));
```

```
for j=2:nx-1
```

```
d2p(j)=(p(j+1)-2*p(j)+p(j-1))/dx^2; % space derivative end
```

```
pnew=2*p-pold+d2p*dt^2; % time extrapolation
pnew(nx/2)=pnew(nx/2)+src(i)*dt^2; % add source term
pold=p; % time levels
p=pnew;
```

```
p(1)=0; % set boundaries pressure free
p(nx)=0;
```

```
% Display
plot(x,p,'b-')
title(' FD ')
drawnow
```

 end



Snapshot Example







Seismogram Dispersion









- Conceptually the most simple of the numerical methods and can be learned quite quickly
- Depending on the physical problem FD methods are conditionally stable (relation between time and space increment)
- FD methods have difficulties concerning the accurate implementation of boundary conditions (e.g. free surfaces, absorbing boundaries)
- FD methods are usually explicit and therefore very easy to implement and efficient on parallel computers
- FD methods work best on regular, rectangular grids