


# Introduction to and State-of-the-Art in Earthquake Source Inversion

– Part 1 –

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## Roadmap

### Earthquake Source Inversion

- (1) Introduction & Theory
- (2) Applications & Implications
- (3) Challenges, Developments, Opportunities

- **How to reach me:**
  - [martin.mai@kaust.edu.sa](mailto:martin.mai@kaust.edu.sa)
- **Material**
  - Slides will be made available
  - General theory based on Aki & Richards (2002), or similar text books (Udias et al., 2014; Stein and Wysession (2002) ....)
  - Examples and applications drawn from numerous research papers



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2

## Roadmap

### Earthquake Source Inversion

#### (1) Introduction & Theory

- A brief overview
- Fundamentals
- From point-source to extended-fault modeling

#### (2) Applications & Implications

- Early developments & case studies
- What can we extract from them?
- What to learn from finite-fault source models?

#### (3) Challenges, Developments, Opportunities

- Imaging versus inversion, or combination of both?
- Alternative methods
- Uncertainty quantification



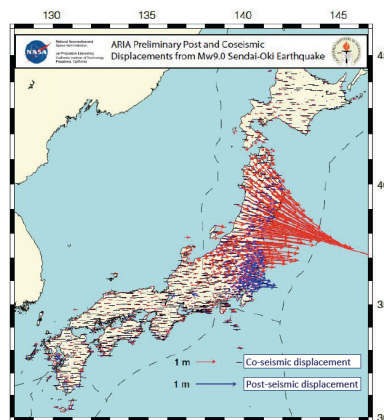
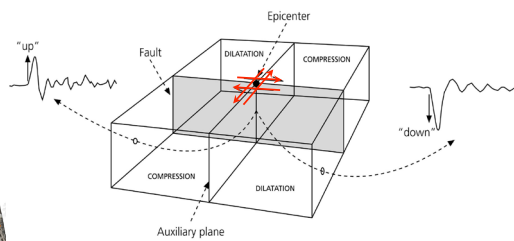
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3

## Introduction & Theory

### A brief overview


- ⊙ Earthquakes occur – mostly – as shear failure on existing faults
- ⊙ As the earthquake happens, it radiates seismic waves that we record and interpret
- ⊙ If the earthquake is large enough, it may generate a measurable permanent offset the Earth surface that can be recorded with satellite-based data (GPS, InSAR)



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4

## Introduction & Theory




### The first earthquake parameters we want to know

- ⊙ Where and when did the earthquake occur?
  - Earthquake location in space and time
    - Numerous algorithms available to solve this weakly non-linear inverse problem (linearized; non-linear; MC ...)
- ⊙ How large was the earthquake?
  - Magnitude estimation
    - Wave-amplitude based rapid estimation
    - Refined estimation requires detailed data analysis
- ⊙ What were the consequences of the earthquake?
  - Damage / impact estimation
    - Macroseismic intensity
    - Engineering assessment

seconds


minutes

days/weeks



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## Introduction & Theory




### Assume size and location are (approximately) known

- ⊙ What happened exactly?
  - Determine which fault(s) ruptured
    - Fault-plane solution
    - Moment-tensor determination
  - Determine the dimensions of the rupture!
    - Aftershock data
    - GPS / InSAR data
    - Waveform analysis
  - Determine the space-time evolution of the rupture!
    - Kinematic finite-fault modelling/inversion
    - Dynamic rupture process

seconds

days

months



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## Introduction & Theory



### Kinematic versus dynamic earthquake source models

#### Kinematic Earthquake Rupture Models

Characterized by time-dependent displacement field (slip vectors) on the rupture plane *without* considering the forces/stresses that cause the motions. The rupture process is entirely specified by the spatio-temporal distribution of the slip direction and the slip-rate function, and how rupture propagates over the fault plane.

#### Dynamic Earthquake Rupture Models

*Build a physical understanding* of the earthquake rupture based on the material properties in the source volume, and the initial and boundary conditions for the forces/stresses acting on the fault plane. The slip-rate vector is obtained by solving the elasto-dynamic equations of motion under the assumption of some constitutive law.

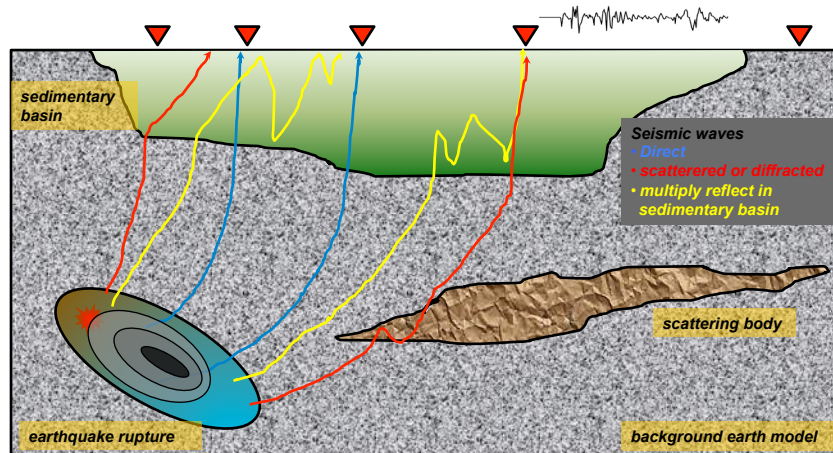


## Introduction & Theory



### Seismic source & wave propagation = linear filter

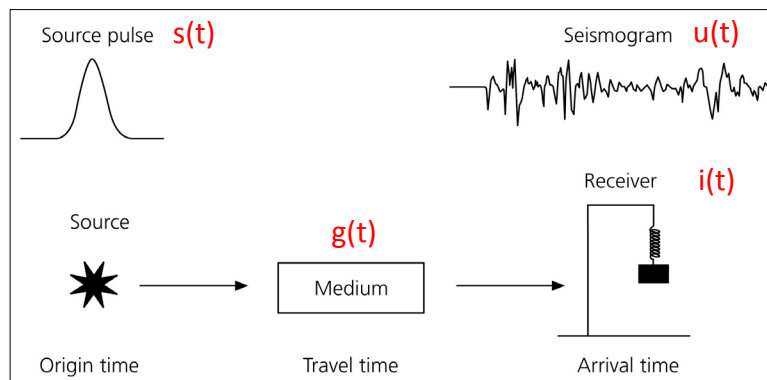
- Conceptual diagram of waves emitted from an elliptical source



## Introduction & Theory

### Seismic source & wave propagation = linear filter

- Conceptual diagram of seismic wave generation & recording



adapted after Stein & Wysession, 2002

## Introduction & Theory

### Seismic source & wave = linear filter

- Generation & propagation & recording of seismic wave as a linear filter

$$u_k(t) = s(t) * g_k(t) * i_k(t)$$

- In this notation:
  - $u(t)$ : observable ground displacement (or velocity) (~known)
  - $s(t)$ : source term (unknown)
  - $g(t)$ : response of the medium (Green's function) (~unknown)
  - $i(t)$ : instrument response (known)

- In the frequency domain

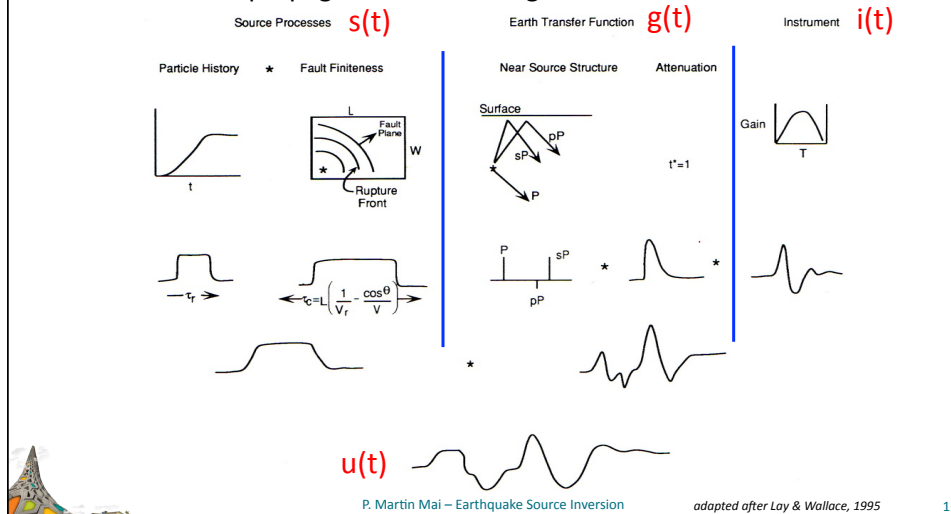
$$U_k(\omega) = S(\omega) \cdot G_k(\omega) \cdot I_k(\omega)$$

## Introduction & Theory



### Seismic source & wave = linear filter

- Generation & propagation & recording of seismic wave as a linear filter



## Introduction & Theory



### For earthquake source modeling / inversion

- Instrument response is (generally) known
  - Could be problematic when older earthquakes are studied
- How to get the Earth' transfer function (Green's function)?
  - Observational
    - "empirical" Green's function, from small nearby events
  - Theoretical / Numerical
    - Calculation for assumed/known Earth model
- How do we want to represent the source term  $s(t)$ ?
  - Point-source
  - Extended-fault source



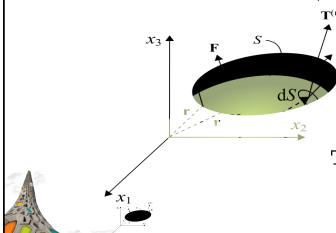
## Introduction & Theory

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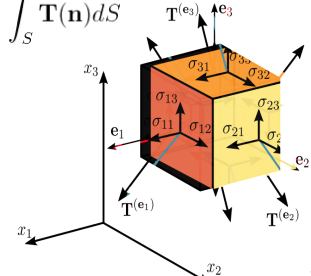
### From equations of motion to Volterra's Theorem

- The particle displacements  $\mathbf{u} = \mathbf{u}(\underline{x}, t)$  (in the Earth or at its surface) are derived as function of space and time in the volume  $V$  (with surface  $S$ ) due to forces applied within and at the surface of that volume
- Assume small deformations (infinitesimal strain); use tractions  $\mathbf{T}$  and stresses  $\boldsymbol{\tau}$  (or  $\boldsymbol{\sigma}$ ) to analyze forces acting between adjacent particles in the volume.
- The momentum equation (" $F = ma$ ") places constraints on accelerations, body forces and tractions acting throughout the volume  $V$  with surface  $S$ .)

$$\frac{\partial}{\partial t} \iiint_V \rho \frac{\partial \mathbf{u}}{\partial t} dV = \iiint_V \mathbf{f} dV + \iint_S \mathbf{T}(\mathbf{n}) dS$$



$\mathbf{T}(\mathbf{n}) = \lim_{dS \rightarrow 0} \frac{\mathbf{F}}{dS}$



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### From equations of motion to Volterra's Theorem

- Given the stress tensor,  $\boldsymbol{\tau}_{kl}$ , the tractions on the infinitesimal small tetrahedron are related to the corresponding normal vectors

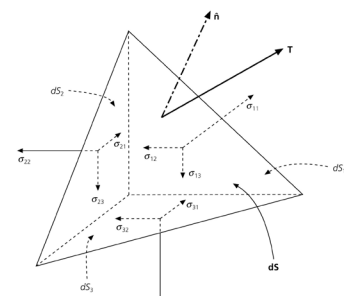
$$\boldsymbol{\tau}_{kl} = \boldsymbol{\sigma}_{kl} = T_i(\hat{x}_k)$$

- $\boldsymbol{\tau}_{kl}$  is the  $l^{\text{th}}$  component of the traction acting across the plane normal to the  $k^{\text{th}}$  axis due to material with larger  $x_k$  acting upon material with smaller  $x_k$ ; the stress-tensor is symmetric

$$T_i = \tau_{ji} n_j$$

- Equations of motion, simplified

$$\rho \ddot{u}_i = f_i + \tau_{ij,j}$$



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## Introduction & Theory



### From equations of motion to Volterra's Theorem

- ⊙ Making a long story short ....

- Include strain tensor  
*(internal deformation; no rigid body motion or rotation)*

$$\varepsilon_{kl} = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

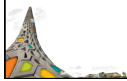
- Write stress-strain relation

$$\tau_{ij} = c_{ijkl} \varepsilon_{kl}$$

- For isotropic, linear elastic media

$$c_{ijkl} = \lambda \delta_{ij} \delta_{kl} + \mu (\delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk})$$

$$\tau_{ij} = \lambda \varepsilon_{kk} \delta_{ij} + 2\mu \varepsilon_{ij}$$



## Introduction & Theory



### From equations of motion to Volterra's Theorem

- ⊙ Making a long story short ....

- To parameterize seismic source we need additional ingredients, i.e. mathematical relations or theorems that connect the displacement  $\mathbf{u}(\mathbf{x}, t)$  to initial conditions and forces
- **Uniqueness Theorem:** *The displacement field  $\mathbf{u} = \mathbf{u}(\mathbf{x}, t)$  in the volume  $V$  with surface  $S$  is **uniquely** defined after time  $t_0$  by initial values of displacement and particle velocity in  $V$ , and by values of the body forces  $\mathbf{f}$ , the tractions  $\mathbf{T}$  over any part  $S_1$  of  $S$  for times  $t \geq t_0$ , and the displacements over the remainder  $S_2$  of  $S$  ( $S_1 + S_2 = S$ )*
- **Betti's theorem:** *Assume displacement field  $\mathbf{u} = \mathbf{u}(\mathbf{x}, t)$  due to body forces  $\mathbf{f}$  and boundary conditions on  $S$  and initial conditions at time  $t = 0$ . Assume further field  $\mathbf{v} = \mathbf{v}(\mathbf{x}, t)$  due to body forces  $\mathbf{g}$ . Using the notation  $\mathbf{T}(\mathbf{u}, \mathbf{n})$  and  $\mathbf{T}(\mathbf{v}, \mathbf{n})$  for the tractions due to  $\mathbf{u}$  and  $\mathbf{v}$ , respectively, the following reciprocal relation is given*

$$\begin{aligned} & \int \int \int_V (\mathbf{f} - \rho \ddot{\mathbf{u}}) \cdot \mathbf{v} dV + \int \int_S \mathbf{T}(\mathbf{u}, \mathbf{n}) \cdot \mathbf{v} dS \\ &= \int \int \int_V (\mathbf{g} - \rho \ddot{\mathbf{v}}) \cdot \mathbf{u} dV + \int \int_S \mathbf{T}(\mathbf{v}, \mathbf{n}) \cdot \mathbf{u} \end{aligned}$$





## Introduction & Theory



### From equations of motion to Volterra's Theorem

⊙ Making a long story short ....

- Betti's theorem does not involve any initial condition for  $\mathbf{u}$  and  $\mathbf{v}$ , and remains true if all related quantities (particle velocities, tractions, body forces) are evaluated at two different times
- Betti's theorem relates motions in the source region to some (far-)distant (surface) displacements, recorded as seismic waves, independent of when the corresponding forces have been applied
- "Trick" : replace the general term  $\mathbf{v}(\mathbf{x}, t)$  in Betti's theorem with a "simple" system response, i.e. the Green's function due to a unidirectional unit pulse precisely located in space and time, and acting in the  $n$ -direction

$$G_{in}(\underline{x}, t; \underline{\xi}, \tau) = A \delta(\underline{x} - \underline{\xi}) \delta(t - \tau) \delta_{in}$$

- B.C's can be chosen such that  $G(\underline{x}, t; \underline{\xi}, \tau)$  has spatial & temporal reciprocity

$$G_{in}(\underline{x}, t; \underline{\xi}, \tau) = G_{ni}(\underline{\xi}, t; \underline{x}, \tau) \quad G_{in}(\underline{x}, t - \tau; \underline{\xi}, 0)$$



## Introduction & Theory



### From equations of motion to Volterra's Theorem

⊙ Making a long story short ....

- Tractions, Green's function, Betti's theorem, Uniqueness

$$u_n(\mathbf{x}, t) = \int_{-\infty}^{\infty} d\tau \int \int \int_V f_i(\xi, \tau) \cdot G_{in}(\xi, t - \tau; \mathbf{x}, 0) dV +$$

volume forces

$$\int_{-\infty}^{\infty} d\tau \int \int_S [G_{in}(\xi, t - \tau; \mathbf{x}, 0) \cdot T_i(\mathbf{u}(\xi, \tau), \mathbf{n})] dS -$$


stress glut

$$\int_{-\infty}^{\infty} d\tau \int \int_S [u_i(\xi, \tau) \cdot c_{ijkl} \cdot n_j \cdot G_{kn,l}(\xi, t - \tau; \mathbf{x}, 0)] dS$$

dislocation



## Introduction & Theory




### Representation Theorem (Volterra's formula)

- ⊙ Neglecting body forces acting on  $\mathbf{u}$ , we can write the representation theorem for internal dislocation sources

$$u_n(\mathbf{x}, t) = \int_{-\infty}^{\infty} d\tau \int_{\Sigma} [u_i(\xi, \tau)] \cdot c_{ijpq} \cdot \nu_j \cdot \frac{\partial}{\partial \xi_q} G_{np}(\mathbf{x}, t - \tau; \xi, 0) d\Sigma$$


$u_n(\mathbf{x}, t)$	n <sup>th</sup> component of the observed displacement field at $\mathbf{x}, t$
$[u_i(\xi, \tau)]$	i <sup>th</sup> component of the displacement at the source as a function of position $\mathbf{x}$ on the fault plane and time $t$
$G_{np}(\mathbf{x}, t - \tau; \xi, 0)$	system response in $n$ -direction due to unit impulse in direction $p$ on the fault plane at $\mathbf{x}, t$
$\frac{\partial}{\partial \xi_q} G_{np}(\mathbf{x}, t - \tau; \xi, 0)$	generalized force couple in $x_q$ -direction and force in $p$ -direction
$c_{ijpq} \cdot \nu_j$	elasticity tensor; normal to the rupture plane



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19

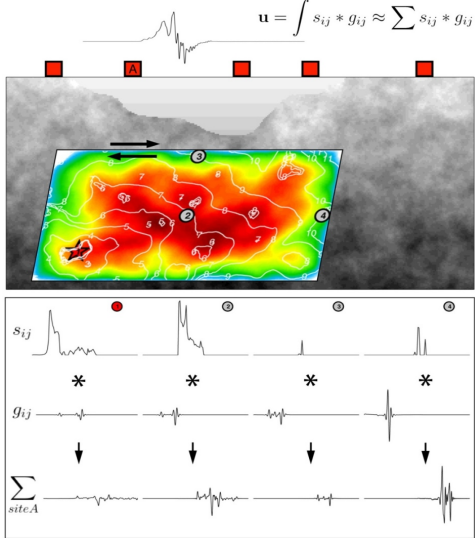
## Introduction & Theory




### Representation Theorem (Volterra's formula)

- ⊙ Conceptually

$$\mathbf{u} = \int s_{ij} * g_{ij} \approx \sum s_{ij} * g_{ij}$$





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20

## Introduction & Theory



### Representation Theorem (Volterra's formula)

- ⊙ In practice
  - The earthquake rupture process, here denoted as  $s_{ij}$ , can be very complex, occurring in a geologically complicated environment. The representation theorem allows us in practice to compute the ground motions
  - The rupture process could be occurring on an extended plane, or multiple planes, but the representation theorem still holds
  - To compute  $u(\underline{x}, t)$  we simply need to sum up all contributions from all points on the fault plane(s), i.e. convolve the local source contribution with the local Green's function (valid for a given site of interest) to obtain a "local" seismogram, and then sum all these local seismograms.
  - **The representation theorem is simply a giant book-keeping device for computing ground motions for arbitrarily complex source models embedded in arbitrarily complex geologic structures.**

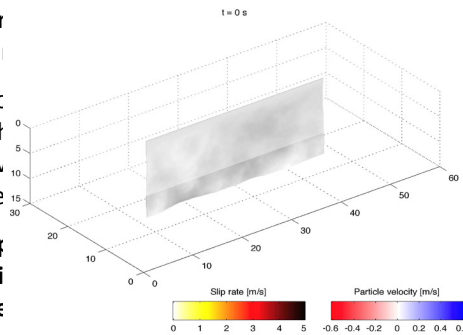


## Introduction & Theory



### Representation Theorem (Volterra's formula)

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## Introduction & Theory

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### Representation Theorem (Volterra's formula)

- ⊙ In practice
  - The earthquake rupture process, here denoted as  $s_{ij}$ , can be very complex,

$t = 0 \text{ s}$

Slip rate [m/s]      Particle velocity [m/s]

23

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### Representation Theorem (Volterra's formula)

- ⊙ In practice
  - The earthquake rupture process, here denoted as  $s_{ij}$ , can be very complex,

24

## Introduction & Theory

### Simplifying the representation theorem

⊙ From complex to point-source

- Space-time heterogeneity of  $[u_i(\xi, \tau)]$ ; equivalent time-dependent point-forces (moment density tensor  $m_{pq}$ )

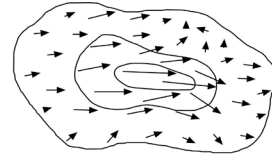
$$u_n(\underline{x}, t) = \iint_{\Sigma} m_{pq} * \frac{\partial}{\partial \xi_q} G_{np} d\Sigma$$

- Earthquake rupture process using space-time averaged source parameters (l: unit vector in direction of slip; k: normal to fracture plane)

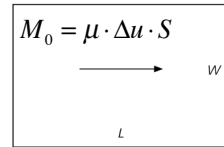
$$u_n(\underline{x}, t) = (l_n k_p + l_p k_n) \int_{-\infty}^{\infty} M_0(\tau) \cdot G_{np,q}(t - \tau) d\tau$$

- Point-source approximation using the point-source double-couple approximation;  $M_{pq}$ : moment tensor

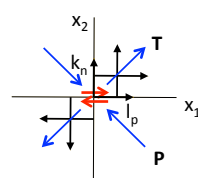
$$u_n(\underline{x}, t) = M_{pq} * G_{np,q}$$



≈



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adapted after Stein & Wysession, 2002

25

## Introduction & Theory

### Point-source shear-fracture (dislocation)

- ⊙ After much algebra, the displacement  $u_n(x,t)$  for a shear-fault with time-dependent moment and moment-rate function  $M_{pq}(t)$  and  $dt M_{pq}(t)$

$$\begin{aligned} M_{pq} * G_{np,q} = & \left( \frac{15\gamma_n \gamma_p \gamma_q - 3\gamma_n \delta_{pq} - 3\gamma_p \delta_{nq} - 3\gamma_q \delta_{np}}{4\pi\rho} \right) \frac{1}{r^4} \int_{r/\alpha}^{r/\beta} \tau M_{pq}(t - \tau) d\tau \\ & + \left( \frac{6\gamma_n \gamma_p \gamma_q - \gamma_n \delta_{pq} - \gamma_p \delta_{nq} - \gamma_q \delta_{np}}{4\pi\rho\alpha^2} \right) \frac{1}{r^2} M_{pq} \left( t - \frac{r}{\alpha} \right) \\ & - \left( \frac{6\gamma_n \gamma_p \gamma_q - \gamma_n \delta_{pq} - \gamma_p \delta_{nq} - 2\gamma_q \delta_{np}}{4\pi\rho\beta^2} \right) \frac{1}{r^2} M_{pq} \left( t - \frac{r}{\beta} \right) \\ & + \frac{\gamma_n \gamma_p \gamma_q}{4\pi\rho\alpha^3} \frac{1}{r} \dot{M}_{pq} \left( t - \frac{r}{\alpha} \right) - \left( \frac{\gamma_n \gamma_p - \delta_{np}}{4\pi\rho\beta^3} \right) \gamma_q \frac{1}{r} \dot{M}_{pq} \left( t - \frac{r}{\beta} \right) \end{aligned}$$

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26

## Introduction & Theory



### Point-source shear-fracture (dislocation)

- After much algebra, the displacement  $u_n(x,t)$  for a shear-fault with time-dependent moment and moment-rate function  $M_{pq}(t)$  and  $dt M_{pq}(t)$

The displacement field due to a general 2<sup>nd</sup>-order moment tensor contains:

- near-field terms, proportional to  $\frac{1}{r^4} \int_{r/\alpha}^{r/\beta} \tau M_{pq}(t - \tau) d\tau$
- Intermediate-field terms,  $\frac{1}{r^2} M_{pq} \left( t - \frac{r}{\alpha} \right)$
- far-field term, proportional to  $\frac{1}{r} \dot{M}_{pq} \left( t - \frac{r}{\beta} \right)$



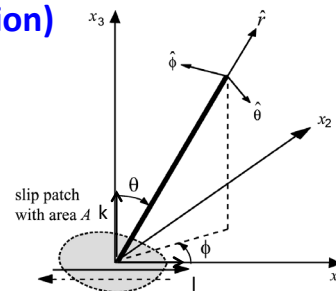
## Introduction & Theory



### Point-source shear-fracture (dislocation)

- In spherical coordinates

$$\begin{aligned}
 \mathbf{u}(x, t) = & \frac{1}{4\pi\rho} A^N \frac{1}{r^4} \int_{r/v_P}^{r/v_S} \tau \dot{M}_0(t - \tau) d\tau \\
 & + \frac{1}{4\pi\rho v_P^2} A^{IP} \frac{1}{r^2} \dot{M}_0(t - r/v_P) \\
 & + \frac{1}{4\pi\rho v_S^2} A^{IS} \frac{1}{r^2} \dot{M}_0(t - r/v_S) \\
 & + \frac{1}{4\pi\rho v_P^3} A^{FP} \frac{1}{r} \dot{M}_0(t - r/v_P) \\
 & + \frac{1}{4\pi\rho v_S^3} A^{FS} \frac{1}{r} \dot{M}_0(t - r/v_S).
 \end{aligned}$$



$$A^N = 9 \sin 2\theta \cos \phi \hat{r} - 6(\cos 2\theta \cos \phi \hat{\theta} - \cos \theta \sin \phi \hat{\phi}),$$

$$A^{IP} = 4 \sin 2\theta \cos \phi \hat{r} - 2(\cos 2\theta \cos \phi \hat{\theta} - \cos \theta \sin \phi \hat{\phi}),$$

$$A^{IS} = -3 \sin 2\theta \cos \phi \hat{r} + 3(\cos 2\theta \cos \phi \hat{\theta} - \cos \theta \sin \phi \hat{\phi}),$$


$$A^{FP} = \sin 2\theta \cos \phi \hat{r},$$

$$A^{FS} = \cos 2\theta \cos \phi \hat{\theta} - \cos \theta \sin \phi \hat{\phi},$$

**SV**                      **SH**



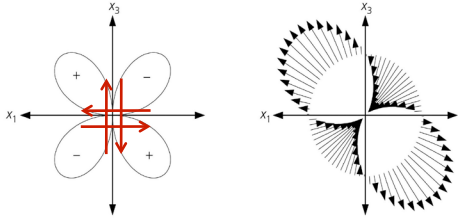
## Introduction & Theory



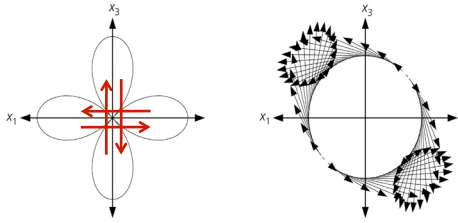
### Point-source shear-fracture (dislocation)

- ⊙ The resulting far-field radiation patterns

#### P-waves




#### S-waves



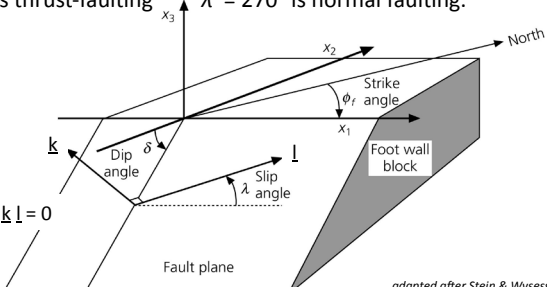
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adapted after Stein & Wyession, 2002
29

## Introduction & Theory




### The fault geometry

- ⊙ Defining geographical parameters of the fault plane
  - Fault dip  $\delta$ , measured from horizontal downward into fault plane
  - Strike  $\phi_f$  measured from north in a RHS coordinate system, hence  $\delta \leq 90^\circ$
  - Rake-angle  $\lambda$  measures the direction of slip in the fault plane as the movement of the hanging wall with respect to the foot wall.
    - $\lambda = 0^\circ$  indicates left-lateral slip,     $\lambda = 180^\circ$  indicates right-lateral slip
    - $\lambda = 90^\circ$  denotes thrust-faulting     $\lambda = 270^\circ$  is normal faulting.




$\mathbf{k} \cdot \mathbf{l} = 0$ 
adapted after Stein & Wyession, 2002
30

## Introduction & Theory




**Assume size and location are (approximately) known**

- What happened exactly?
  - Determine which fault(s) ruptured!
    - Fault-plane solution
    - Moment-tensor determination



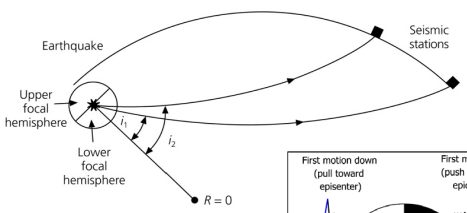
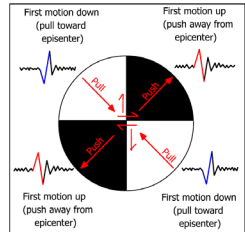
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## Introduction & Theory

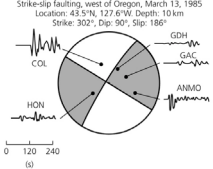


**Assume size and location are (approximately) known**

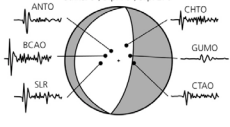
- What happened exactly?
  - Determine which fault(s) ruptured!
    - Fault-plane solution from first motion polarities
    - Lower-hemisphere projection

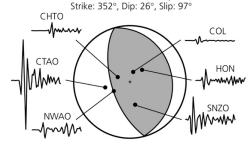
Strike-slip faulting, west of Oregon, March 13, 1985  
Location: 43.5°N, 127.6°W, Depth: 10 km  
Strike: 302°, Dip: 90°, Slip: 186°



Normal faulting, mid-Indian rise, May 16, 1985  
Location: 29.1°S, 77.7°E, Depth: 10 km  
Strike: 8°, Dip: 70°, Slip: 210°



Thrust faulting, Vanuatu Islands, July 3, 1985  
Location: 17.2°S, 167.8°E, Depth: 30 km  
Strike: 352°, Dip: 26°, Slip: 97°



rce Inversion adapted after Stein & Wyssession, 2002

32



## Introduction & Theory

**Assume size and location are (approximately) known**

⊙ What happened exactly?

- Determine which fault(s) ruptured!
  - Moment-tensor determination  $u_n(\underline{x}, t) = M_{pq} * G_{np,q}$

$$M_{pq} = \mathbf{M} = \begin{pmatrix} M_{11} & M_{12} & M_{13} \\ M_{21} & M_{22} & M_{23} \\ M_{31} & M_{32} & M_{33} \end{pmatrix} = \begin{pmatrix} M_{xx} & M_{xy} & M_{xz} \\ M_{yx} & M_{yy} & M_{yz} \\ M_{zx} & M_{zy} & M_{zz} \end{pmatrix} = \begin{pmatrix} M_{\theta\theta} & M_{\theta\phi} & M_{\theta r} \\ M_{\phi\theta} & M_{\phi\phi} & M_{-\phi r} \\ M_{r\theta} & M_{-\phi r} & M_{rr} \end{pmatrix}$$

P. Martin Mai – Earthquake Source Inversion

## Introduction & Theory

**Assume size and location are (approximately) known**

⊙ What happened exactly?

- Determine which fault(s) ruptured!
  - Moment-tensor determination  $u_n(\underline{x}, t) = M_{pq} * G_{np,q}$


$$M_{pq} = \mathbf{M} = \begin{pmatrix} M_{11} & M_{12} & M_{13} \\ M_{21} & M_{22} & M_{23} \\ M_{31} & M_{32} & M_{33} \end{pmatrix} = \begin{pmatrix} M_{xx} & M_{xy} & M_{xz} \\ M_{yx} & M_{yy} & M_{yz} \\ M_{zx} & M_{zy} & M_{zz} \end{pmatrix} = \begin{pmatrix} M_{\theta\theta} & M_{\theta\phi} & M_{\theta r} \\ M_{\phi\theta} & M_{\phi\phi} & M_{-\phi r} \\ M_{r\theta} & M_{-\phi r} & M_{rr} \end{pmatrix}$$

- Isotropic explosion/implosion
- Vertical strike-slip
- Vertical dip-slip
- 45°-dipping thrust
- Two versions of non-double couple mechanisms

Moment tensor	Beachball	Moment tensor	Beachball
$\frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$		$\frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix}$	
$\frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$		$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$	
$\frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 0 & -1 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{pmatrix}$		$\frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & -1 & 0 \end{pmatrix}$	
$\frac{1}{\sqrt{2}} \begin{pmatrix} -1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$		$\frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$	
$\frac{1}{\sqrt{6}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & -2 & 0 \\ 0 & 0 & 1 \end{pmatrix}$		$\frac{1}{\sqrt{6}} \begin{pmatrix} -2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$	
$\frac{1}{\sqrt{6}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$		$\frac{1}{\sqrt{6}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$	

P. Martin Mai – Earthquake Source Inversion adapted after Stein & Wysession, 2002

## Introduction & Theory




**Assume size and location are (approximately) known**

- ⊙ What happened exactly?
  - Determine which fault(s) ruptured!
    - **Moment-tensor inversion**
      - 6 independent components; 5 if only deviatoric ( $M_{11} + M_{22} + M_{33} = 0$ )
      - Express as  $M_{pq}(t) = M_0 \cdot f(t) \cdot m_{pq}$  and assume  $f(t)$

$$u_n(\underline{x}, t) = \sum_{i=1}^6 g_i^n(r, t) * m_i(t) \longrightarrow \underline{u} = \underline{G} \underline{m}$$


- Solution for  $\underline{m}$  depends on the structure of  $\underline{G}$  (the inverse problem / data)
 
$$\underline{m} = (\underline{G}^T \underline{G})^{-1} \underline{G}^T \underline{u} \qquad \underline{m} = (\underline{G}^T \underline{G} + \lambda \underline{R}^T \underline{R})^{-1} \underline{G}^T \underline{u}$$

$$\underline{G} = \underline{U} \underline{\Lambda} \underline{V}^T \qquad \underline{m} = (\underline{V} \underline{\Lambda}^{-1} \underline{U}^T) \underline{u}$$



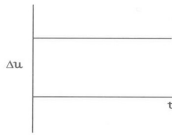
P. Martin Mai – Earthquake Source Inversion 35

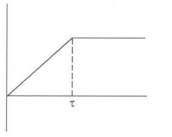
## Introduction & Theory

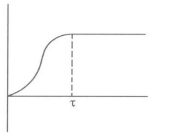


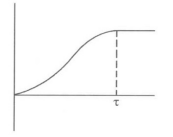
**Assume size and location are (approximately) known**


- ⊙ What happened exactly?
  - Determine which fault(s) ruptured!
    - Moment-tensor determination  $u_n(\underline{x}, t) = M_{pq} * G_{np,q}$
    - Temporal dependence,  $f(t)$   $M_0(t - \tau) = M_{pq}(t - \tau)$

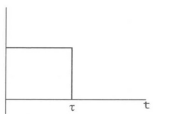


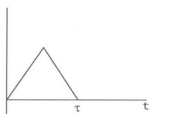


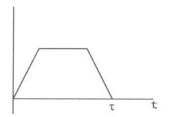













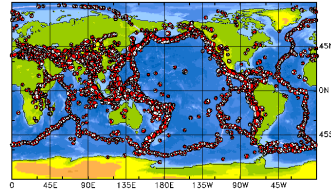


P. Martin Mai – Earthquake Source Inversion 36

# Introduction & Theory

## Some simple cases studies

- ⊙ CMT (Centroid Moment Tensor) Solution
  - Adjust for the fact that the seismic moment is not released at the rupture nucleation point (hypocenter)



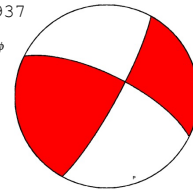
(<http://www.globalcmt.org/>)

### 110302J CENTRAL ALASKA

Date: 2002/11/ 3 Centroid Time: 22:13:28.0 GMT  
 Lat= 63.23 Lon=-144.89  
 Depth= 15.0 Half duration=23.5  
 Centroid time minus hypocenter time: 47.0  
 Moment Tensor: Expo=27 0.513 -6.038 5.525 0.183 2.615 -3.937

$$M_{rr} \quad M_{\theta\theta} \quad M_{\phi\phi} \quad M_{\theta r} \quad M_{\phi r} \quad M_{\theta\phi}$$

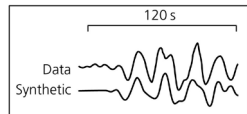
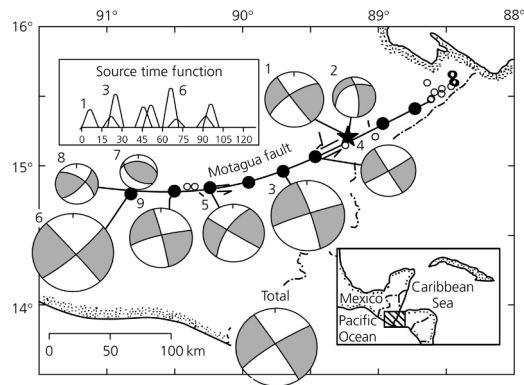
Mw = 7.8 mb = 7.0 Ms = 8.5 Scalar Moment = 7.48e+27  
 Fault plane: strike=296 dip=71 slip=171  
 Fault plane: strike=29 dip=82 slip=19



# Introduction & Theory

## Some simple cases studies


- ⊙ Wave-form modeling
  - Add multiple moment-tensor contributions to fit waveforms for a complex earthquake



*adapted after Stein & Wyession, 2002*




## Introduction & Theory



### Standard finite-fault earthquake source inversion


- The earthquake-source term is space-time dependent, i.e. each “point” on the fault is activated to slip when it is reached by the expanding rupture front
- Rupture speed and the local slip-rate function depend on initial & boundary conditions (stress, friction etc.); both may vary strongly over the rupture plane
- The earthquake rupture process itself is a highly nonlinear process
- However, standard/classical earthquake source inversions approaches “linearize” the problem
  - What exactly do we mean by “linearize” here?
  - Which assumptions do we have to make?



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39

## Introduction & Theory



### Standard finite-fault earthquake source inversion


- The representation theorem is linearized, using a number of assumptions (e.g. Olson & Apsel, 1982; Hartzell & Heaton, 1983):

$$u_n(\mathbf{x}, t) = \int_{-\infty}^{\infty} d\tau \int_{\Sigma} [u_i(\xi, \tau)] \cdot \nu_j \cdot c_{ijpq} \cdot \frac{\partial}{\partial \xi_q} G_{np}(\mathbf{x}, t - \tau; \xi, 0) d\Sigma$$

Seismic data observations <i>what we have</i>	Space-time dependent rupture process <i>what we want</i>	Earth structure and related Green's functions <i>what we need to compute</i>
--------------------------------------------------	-------------------------------------------------------------	---------------------------------------------------------------------------------

$d = g(m)$ 
 $\underline{d} = \underline{G} \underline{m}$

- Assuming we “know” **Earth structure & fault geometry**, we can examine seismic observations to learn about **earthquake source properties** given **seismological and/or geodetic data**.
- This is a kinematic representation; we cannot infer rupture physics here!



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40

## Introduction & Theory



### Standard finite-fault earthquake source inversion

- The representation theorem is “linearized” using a number of assumptions (e.g. Olson & Apsel, 1982; Hartzell & Heaton, 1983):
  - Source geometry is known, and approximated by one or more rectangular planes that are again subdivided into a set of “subfaults” (nodes)
  - An elementary identical slip function acts on each point on the fault
  - The slip-history at each point can be represented by a summation of elementary basis functions, lagged in time (multi-time window)
  - The rupture velocity and is constant



## Introduction & Theory



### Standard finite-fault earthquake source inversion

- The representation theorem is linearized, using a number of assumptions (e.g. Olson & Apsel, 1982; Hartzell & Heaton, 1983):

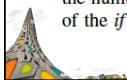
$$u_n(\mathbf{x}, t) = \int_{-\infty}^{\infty} d\tau \int_{\Sigma} [u_i(\xi, \tau)] \cdot \nu_j \cdot c_{ijpq} \cdot \frac{\partial}{\partial \xi_q} G_{np}(\mathbf{x}, t - \tau; \xi, 0) d\Sigma$$

- Discretized version:

$$u_n(\mathbf{x}, t) = \sum_{itm=1}^{ntm} \sum_{is=1}^{ns} \sum_{if=1}^{nf} m(if, is, itm) \times \int [u_{\text{unit}_{is}}(\tau - \Delta t_{\text{trig}})] \\ \times c_{i(is)jkl}(\xi) n_j G_{kn,l}(\mathbf{x}, t - \tau; \xi(if), 0) d\tau$$

Here,  $m(if, is, itm)$  is the amount of slip in the  $is$ th direction at the  $itm$ th time window on the  $if$ th subfault,  $nf$  is the number of subfaults,  $ns$  is the number of slip directions,  $ntm$  is the number of time windows,  $R$  is the hypocentral distance of the  $if$ th subfault, and  $[u_{\text{unit}_{is}}(\tau)]$  is the unit slip function.

$$\Delta t_{\text{trig}} = \frac{R}{V_r} + \Delta tw \cdot (itm - 1)$$



## Introduction & Theory

### Standard finite-fault earthquake source inversion

- Conceptual sketch
  - Geometry assumed or known
  - Rupture speed  $V_r$  assumed
  - Elementary slip (or slip-rate) function assumed

$$\Delta t_{\text{trig}} = \frac{R}{V_r} + \Delta t_w \cdot (itm - 1)$$

△ : hypocenter

From Delouis et al, 2002

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## Introduction & Theory

### Standard finite-fault earthquake source inversion

- The system of equations then looks like

$$\underline{d} = \underline{G}\underline{m} = \sum_j G_{ij} m_j$$

include smoothing

$$\begin{bmatrix} \underline{d} \\ \underline{0} \end{bmatrix} = \begin{bmatrix} \underline{G} \\ \lambda \underline{R} \end{bmatrix} \underline{m}$$


$\underline{R}$  : smoothing (regularization) matrix to account for variations in model parameters with distance and/or time (the farther apart subfaults are, the larger a difference is allowed);  
 $\lambda$  has to be determined by trial-and-error, or some statistical techniques

$$\begin{matrix} \text{(time) data points station 1} \\ \text{(time) data points station 2} \\ \vdots \\ \text{(time) data points station N} \end{matrix} \begin{bmatrix} d_1 \\ d_2 \\ \vdots \\ d_1 \\ d_2 \\ \vdots \\ d_1 \\ d_2 \\ \vdots \end{bmatrix} = \begin{matrix} \text{Subfault 1} \\ \text{Subfault 2} \\ \vdots \\ \text{Subfault m} \end{matrix} \begin{bmatrix} G_{11} & G_{12} & \cdots & G_{1m} \\ G_{21} & G_{22} & \cdots & G_{2m} \\ \vdots & \vdots & & \vdots \\ G_{11} & G_{12} & \cdots & G_{1m} \\ G_{21} & G_{22} & \cdots & G_{2m} \\ \vdots & \vdots & & \vdots \\ G_{11} & G_{12} & \cdots & G_{1m} \\ G_{21} & G_{22} & \cdots & G_{2m} \\ \vdots & \vdots & & \vdots \end{bmatrix} \begin{bmatrix} s_1 \\ s_2 \\ \vdots \\ s_m \end{bmatrix}$$

Dislocation in subfault 1  
Dislocation in subfault 2  
Dislocation in subfault m

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## Introduction & Theory



**Standard finite-fault earthquake source inversion**


- Solution strategies (least-squares)
  - Include data covariance,  $\sigma^2 C_d$ 

$$\underline{m} = (G^T C_d^{-1} G)^{-1} G^T \underline{d}$$
  - Covariance of model parameters
 
$$C_m = \sigma^2 (G^T C_d^{-1} G)^{-1}$$
  - Difficult to use; need constraints (non-negativity; smoothness)
 
$$\underline{m} = (G^T C_d^{-1} G + \lambda^2 R^T R)^{-1} G^T \underline{d} \quad C_m = \sigma^2 (G^T C_d^{-1} G + \lambda^2 R^T R)^{-1}$$

*$\sigma^2$  and  $\lambda^2$  may be chosen by Akaike-Bayesian Information Criterion (ABIC)*

$$ABIC(\lambda) = N_d \ln Res(\underline{m}) - N_m \ln \lambda^2 + \ln |C_m| + b \quad ABIC = -2(\log \text{marginal likelihood}) + 2(\text{number of hyperparameters})$$
- Apply singular-value decomposition to  $G$ 


$$G = U \Lambda V^T \quad \underline{m} = V_s \Lambda_s^{-1} U_s^T \underline{d}$$



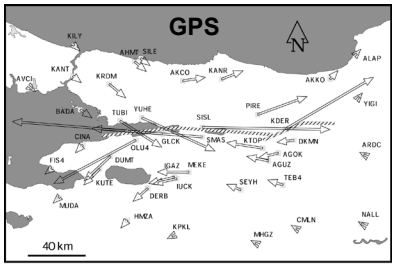
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45

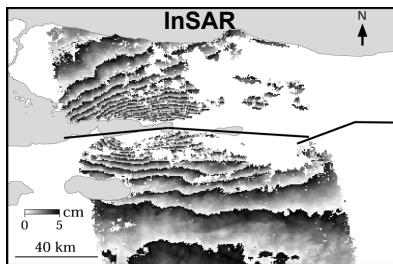
## Introduction & Theory



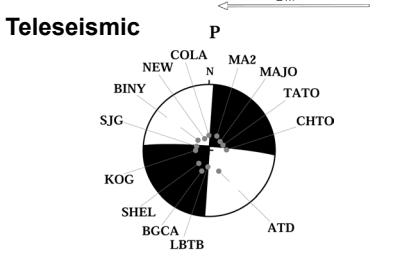
**Example: 1999 Izmit (M 7.6) earthquake**



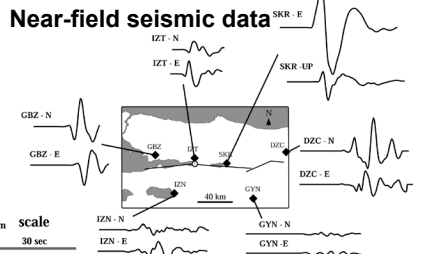
**GPS**




**InSAR**



**Teleseismic**



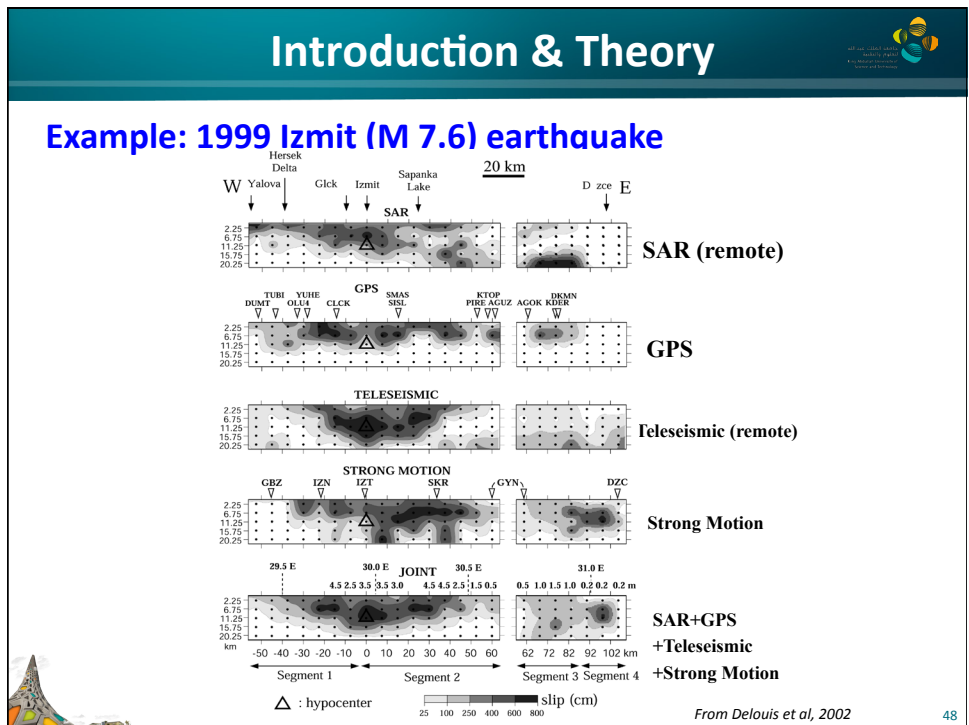
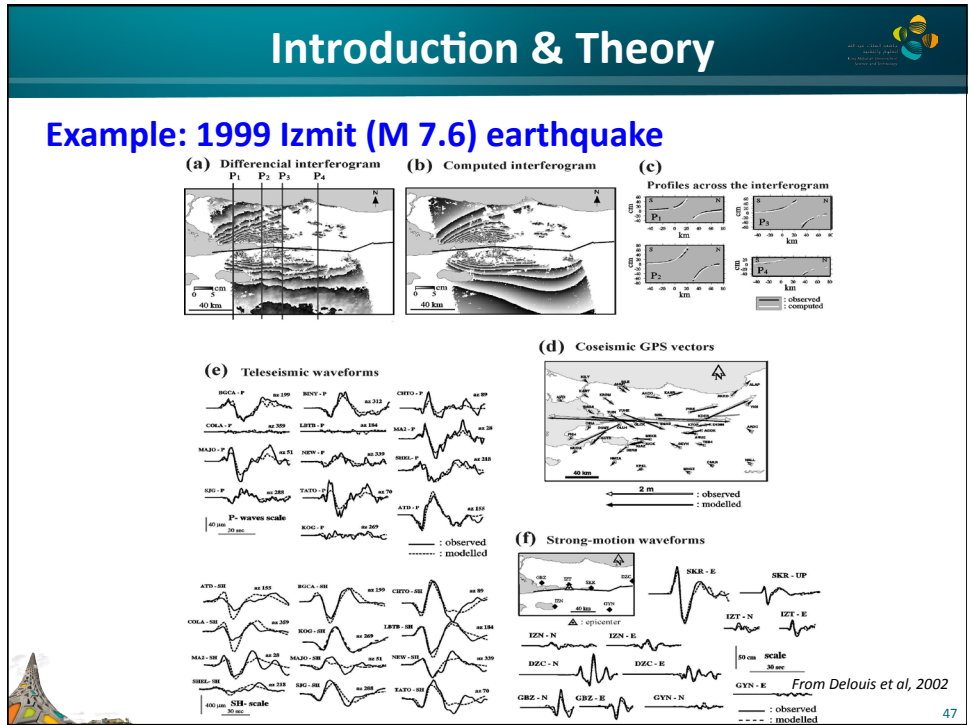
**Near-field seismic data**



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From Delouis et al, 2002

46



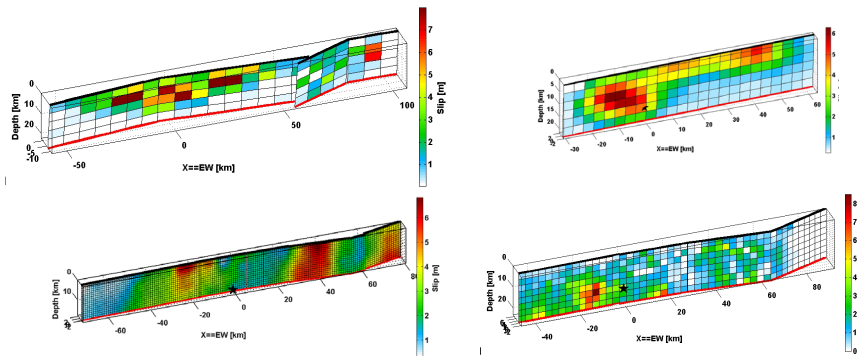


## Introduction & Theory



### Example: 1999 Izmit (M 7.6) earthquake

- Four solutions as available on the SRCMOD database (equake-rc.info/SRCMOD)



From Delouis et al, 2002

49

## Introduction & Theory



### Using permanent (static) surface displacements

- Geodetic data (GPS/InSAR) for source inversions
  - Use analytical Okada (1989; 1992) solutions (elastic half-space) to compute the response of the medium
  - Non-linear inversion for fault geometry (with constant fault): strike, dip, rake, depth, width, length, position (x,y)
  - Once the geometry is fixed, the inversion is linear to find distributed slip on the fault
- If possible, different data sets are combined
  - Geodetic data to constrain geometry and shallow slip
  - Teleseismic to constrain seismic moment and overall characteristics
  - Strong motion data to constrain temporal rupture evolution
  - How to choose the weights for each data set?

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50

## Introduction & Theory

**Using permanent (static) surface displacements**

- Geodetic data (GPS/InSAR) for source inversions: **Hector Mine, 1999**

**Fault Geometry & Data**

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## Introduction & Theory

**Using permanent (static) surface displacements**

- Geodetic data (GPS/InSAR) for source inversions: **Hector Mine, 1999**

**“optimal” Laplacian smoothing (R) and data-dependent final solution**

$$\underline{m} = (G^T G + \lambda^2 R^T R)^{-1} G^T d$$

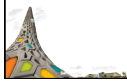
after Jonsson et al, 2002 P. Martin Mai – Earthquake Source Inversion 52

## Introduction & Theory



### Where are the problems? (I)

- ⊙ Green's functions
  - Do we know Earth structure well enough?
  - Do we compute the Green's function correctly?
  - Should we include a formal error term for the Green's function?
  
- ⊙ Fault geometry
  - Unknown a priori, often constrained from aftershocks and – if available – surface-faulting information
  - Slight variations in the geometry will affect the Green's function and hence the entire inversion
  - Different parameterizations of the fault plane

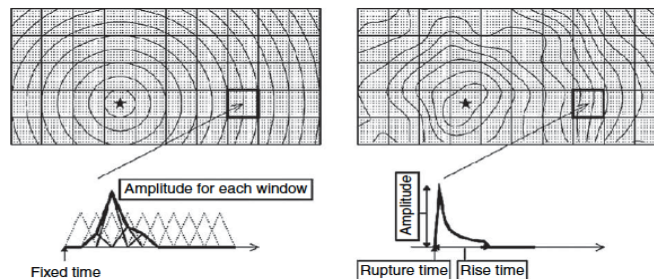


## Introduction & Theory



### Where are the problems? (II)

- ⊙ Parameterization
  - Slip function on each point? How to be more physical?
  - Relaxing the rupture-speed assumption

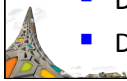


## Introduction & Theory



### Where are the problems? (III)

- ⊙ Solution to the inverse problem
  - Various approaches to solve the linearized problem
    - Optimal choice of smoothing / regularization?
  - Many methods to do a non-linear inversion
    - Building in physical constraints into the search space
  - Formal uncertainty quantification of the resulting model parameters
  - Parameterization of the inversion
  
- ⊙ Data
  - What is the data distribution? Is that sufficient, or biased somehow?
  - Data selection / trimming / processing
  - Data weighting
  - Data uncertainties



## Introduction & Theory



### Next ?

#### (2) Applications & Implications

- Early developments & case studies?
- What can we extract from them?
- What to learn from finite-fault source models?

